

# Overview of the experimental anomalies

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## Abstract

The phenomenon of neutrino oscillations has been established by a number of experiments. An incomplete list includes Homestake, Gallex, Sage, Kamiokande, IMB, Soudan2, Macro, SNO and recently also Double Chooze, Reno and Daya Bay. Assuming existence of three neutrinos, one can extract neutrino mixing angles and mass differences from the data and use them to predict outcome of any other oscillation experiment. However, there is a handful of experiments, whose results are incompatible with this standard picture and hint towards the existence of an additional, sterile, neutrino. Among these are LSND, KARMEN, MiniBooNE, MINOS and many of the short-baseline reactor experiments. In this talk I will review the experimental setups and results, as well as their implications for the neutrino physics.

## Outline

1. **Conventional picture**
2. **Experimental anomalies**
  - a. LSND and KARMEN
  - b. MiniBooNE
  - c. Gallium anomaly
  - d. Reactor anomaly
  - e. MINOS
  - f. Global fits
3. **Summary**

## Conventional picture

- In the standard model neutrinos are strictly massless. Historically it is explained by the fact, that when the Standard Model was formulated, all the existing experiments were compatible with zero neutrino mass and absence of the right-handed neutrino.
- The situation was changed by the **Homestake** experiment, which operated from 1970 till 1994. This experiment has observed only **1/3 of the flux** predicted by the Standard Solar Model. To explain this result the neutrino oscillation hypothesis was proposed. According to this hypothesis the electron neutrinos produced in the sun oscillate to the other flavours on the way to Earth. And since the detector is only capable of measuring the electron,



but not the muon or tau - neutrinos, this explains results of the Homestake experiment.

- A decisive continuation of this hypothe-

- A decisive confirmation of this hypothesis was provided by the **SNO** experiment. It measured **neutral-current** interactions:

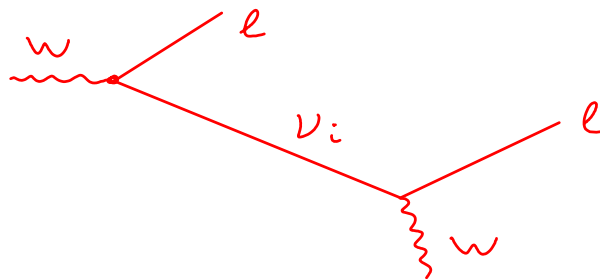


which are equally sensitive to all neutrino flavors. The measured total neutrino flux was in agreement with the prediction of the standard solar model. On the other hand the flux of electron neutrinos measured via **charged current** interaction:



was smaller than the total flux. This has proven that the oscillation hypothesis is correct and that neutrinos have masses.

- A typical oscillation experiment is described by the diagram:



The neutrino is produced as flavor eigenstate but propagates as a mass eigenstate. If the neutrinos are masses, then the

If the neutrinos have masses, then the mass eigenstates coincide with the flavor eigenstates and oscillations are impossible. This implies that the Standard Model should be extended, for instance in the one of the ways discussed by Lisa.

- The oscillation probability depends on the neutrino masses and mixing angles.

$$P_{\alpha \rightarrow \beta} = \left| \sum_i U_{\alpha i} e^{+im_i^2 L/2E} U_{i\beta}^* \right|^2$$

$\alpha \rightarrow i \rightarrow \beta$

We can simplify this expression to:

$$P_{\alpha \rightarrow \beta} = \delta_{\alpha\beta} - 4 \sum_{i>j} \text{Re}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin^2\left(\frac{\Delta m_{ij}^2 L}{4E}\right) + 2 \sum_{i>j} \text{Im}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin\left(\frac{\Delta m_{ij}^2 L}{2E}\right),$$

As follows from this expression, amplitude of the oscillations depends on the mixing angles, and the oscillating behavior on the product  $\Delta m_{ij}^2 (L/E)$ . This dependence is decisive for the whole following discussion.

Changing the ratio  $(L/E)$  we can test different values of  $\Delta m_{ij}^2$ . Loosely speaking, the shorter the distance, the larger is  $\Delta m^2$  one can study.

- The oscillation parameters have been measured by a number of experiments.

measured by a number of experiments.

The **solar** neutrino angle and the corresponding mass difference was measured by **Callex**, **Sage**, **Kamionande**, and lately by **Borexino** experiments.

$$\theta_{12} \approx 33^\circ, \Delta M_{21}^2 \approx 7.5 \times 10^{-5} \text{ eV}^2$$

The **atmospheric** parameters have been measured by **Kamionande**, **IMB**, **Soudan 2**, and **Macro** experiments.

$$\theta_{23} \approx 40^\circ \text{ or } 50^\circ, \Delta M_{31}^2 = 2.47 \times 10^{-3} \text{ eV}^2$$

← assuming normal hierarchy.

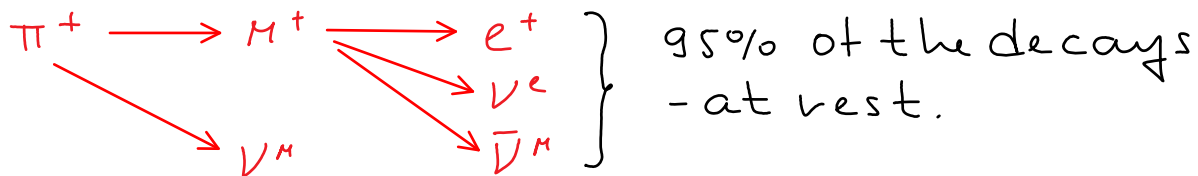
Finally the missing <sup>reactor</sup> mixing angle has recently been measured by **Double Chooz**, **Reno**, and **Daya Bay** experiments.

$$\theta_{13} \approx 8.6^\circ$$

- There also CP-violating phases, one or three, but they are not known at present. But the two mass differences and the three mixing angles are sufficient to predict the outcome of pretty much every neutrino oscillation experiment. A deviation from the predicted behavior is an anomaly.

## LSND, KARMEN

- Historically, one of the first anomalies was found by the **LSND** experiment.
- Details of the experiment:
  - used **800 MeV proton beam** to create pions.
  - neutrinos were produced in the following **decay chain**:



$\pi^-$  - absorbed in the target and the expected  $\bar{\nu}_e$  flux is  $\sim 8 \times 10^{-4}$  of the  $\bar{\nu}_\mu$  flux.

- Primary goal - search for  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  oscillations, where  $\bar{\nu}_e$  are identified by  $\bar{\nu}_e p \rightarrow e^+ n$ ; then  $n + p \rightarrow d + \gamma$  (2.2 MeV)

Physical potential.

$$L = 30 \text{ m}$$

$$E_{\bar{\nu}_\mu} = 20 \dots 60 \text{ MeV}$$

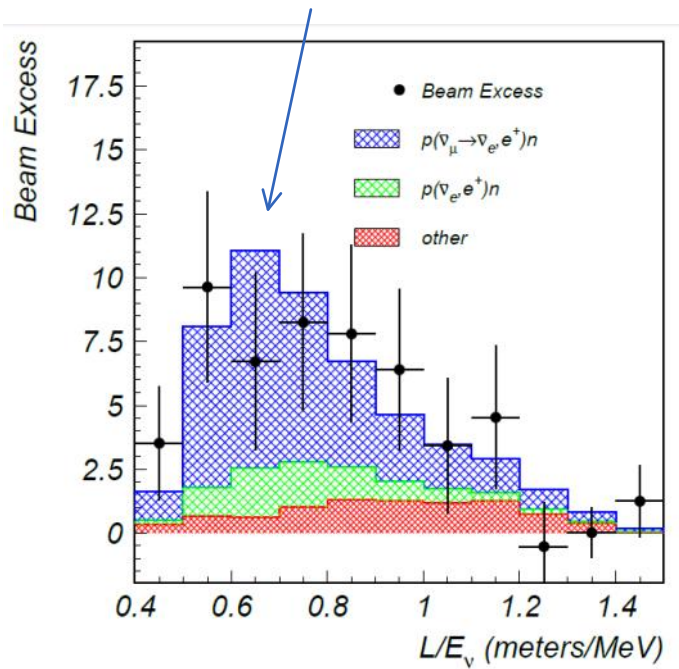
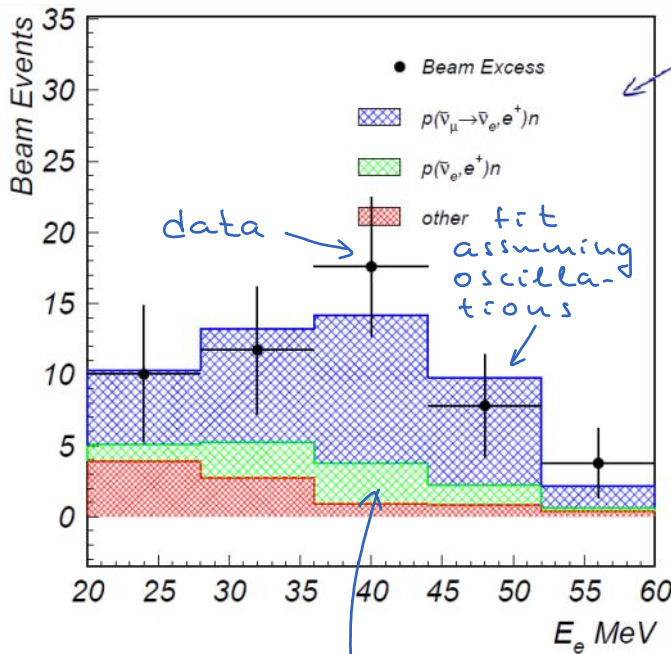
$$F_{\nu} = 60 \dots 200 \text{ M.e.V}$$

$E_{\nu} - 60 - 200 \text{ MeV}$

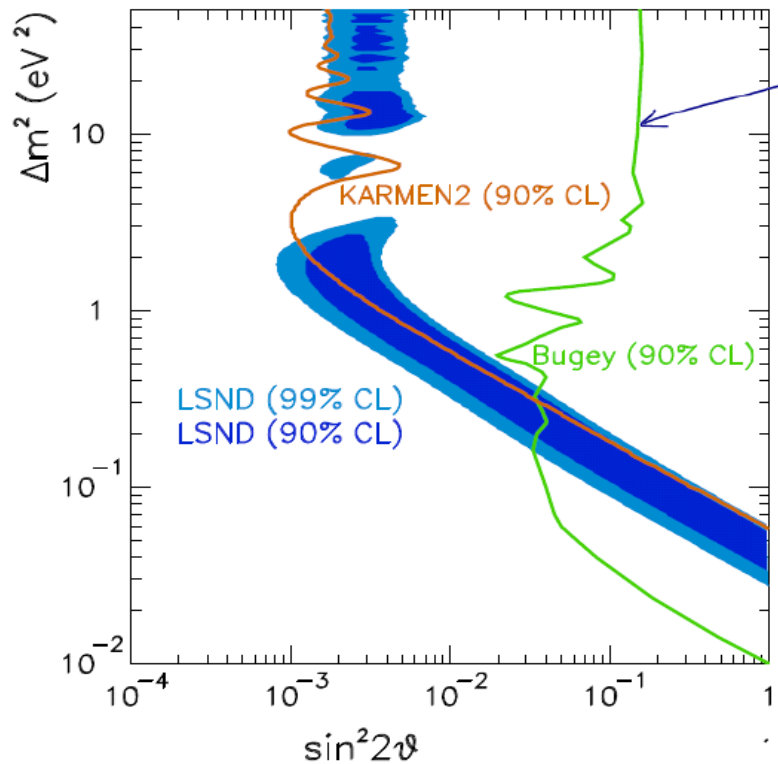
$$\frac{E(1 \text{ MeV})}{L(100 \text{ km})} \sim 10^{-5} \text{ eV}^2 \rightarrow \frac{E}{L} \sim 1 \text{ eV}^2$$

If the conventional picture is correct, neutrinos would not have time to oscillate on such a short distance.

- They have **observed an excess**. Energy and  $L/E$  dependence of the excess:  
*non-trivial L-dependence*



- The data can be fit assuming existence of an additional neutrino with  $\Delta m^2 \sim 1$ . Larger  $\Delta m^2$  are also acceptable.



- A very similar experimental setup was used in the **KARMEN** experiment. The main differences, as compared to LSND are:
  - smaller distance to the detector,  $L \sim 18 \text{ m}$
  - clear separation of  $\nu_\mu$  from  $\pi^+$  decay from  $\bar{\nu}_\mu$  and  $\nu_e$  from  $\mu^+$  decay.

The expected **background** was **15.8** events and the **measured** number was **15**. In other words, KARMEN does not support the claim of LSND. To be fair, one should note that KARMEN does not exclude part of the LSND plot with  $\Delta m^2 < 2 \text{ eV}^2$  and  $\Delta m^2 \sim 7 \text{ eV}^2$ .



# MiniBooNE

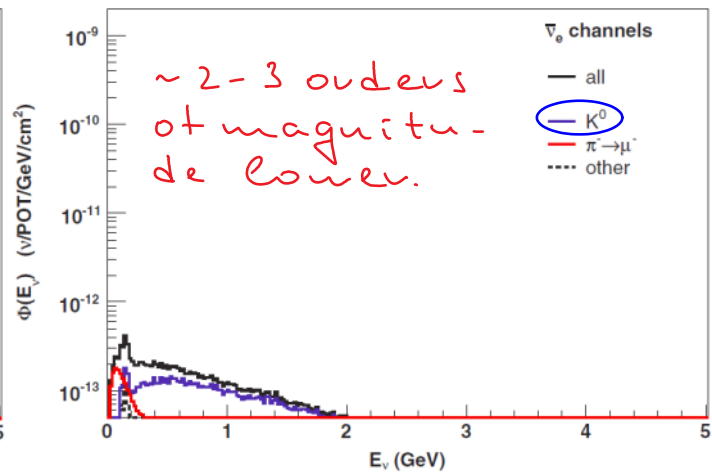
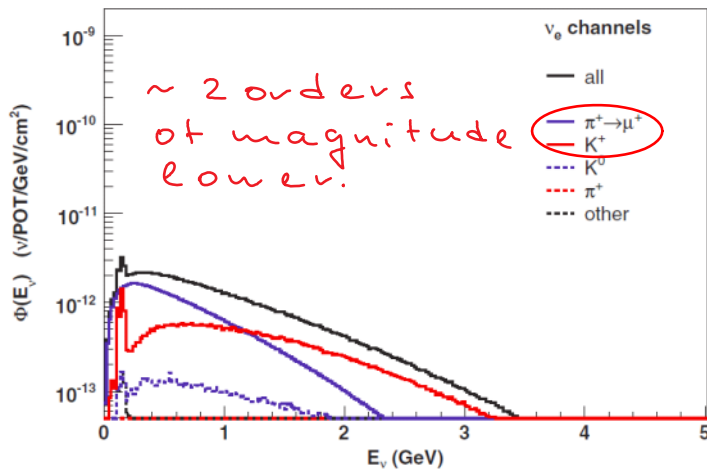
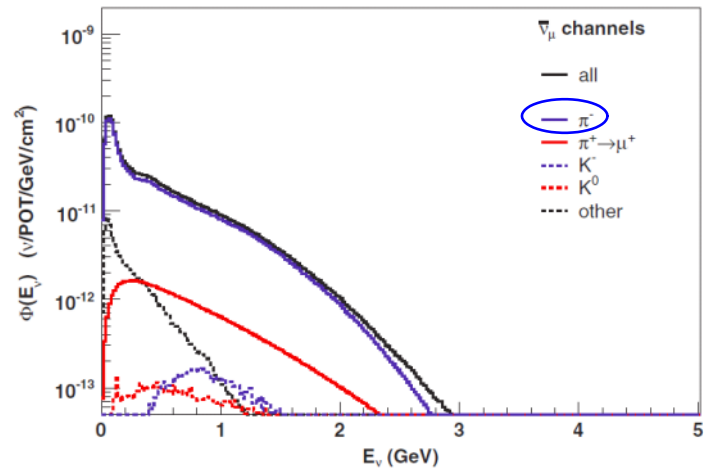
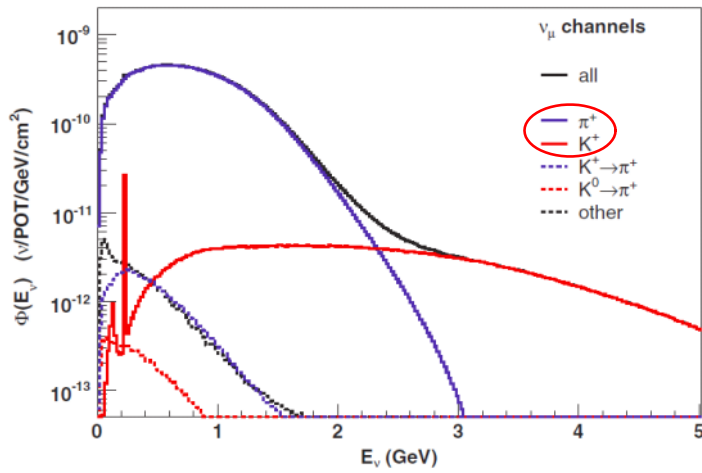
- To cross-check the controversial results of LSND the same group of people has built another experiment - **MiniBooNE**.

Details of the experiment:

- used **8 GeV proton beam** to create pions and kaons.
- magnetic field was used to form either a  **$\pi^+$ ,  $K^+$  or  $\pi^-$ ,  $K^-$  beam**.
- **neutrino / antineutrino modes**

| Particle | Lifetime (ns) | Decay mode                      | Branching ratio (%) |
|----------|---------------|---------------------------------|---------------------|
| $\pi^+$  | 26.03         | $\mu^+ + \nu_\mu$               | 99.9877             |
|          |               | $e^+ + \nu_e$                   | 0.0123              |
| $K^+$    | 12.385        | $\mu^+ + \nu_\mu$               | 63.44               |
|          |               | $\pi^0 + e^+ + \nu_e$           | 4.98                |
|          |               | $\pi^0 + \mu^+ + \nu_\mu$       | 3.32                |
| $K_L^0$  | 51.6          | $\pi^- + e^+ + \nu_e$           | 20.333              |
|          |               | $\pi^+ + e^- + \bar{\nu}_e$     | 20.197              |
|          |               | $\pi^- + \mu^+ + \nu_\mu$       | 13.551              |
|          |               | $\pi^+ + \mu^- + \bar{\nu}_\mu$ | 13.469              |
| $\mu^+$  | 2197.03       | $e^+ + \nu_e + \bar{\nu}_\mu$   | 100.0               |

- neutrino and antineutrino fluxes:



- Primary goal: study of  $\nu_\mu \rightarrow \nu_e$  and  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  oscillations. Thus, it can cross-check the LSND result.
- detection process - CC quasi-elastic scattering:  $\nu_e C \rightarrow e^- X$  and  $\bar{\nu}_e C \rightarrow e^+ X$
- the detector cannot distinguish between  $\nu_e$  and  $\bar{\nu}_e$  candidates
- Physical potential:
  - the detector is located 541 m away from the decay pipe.

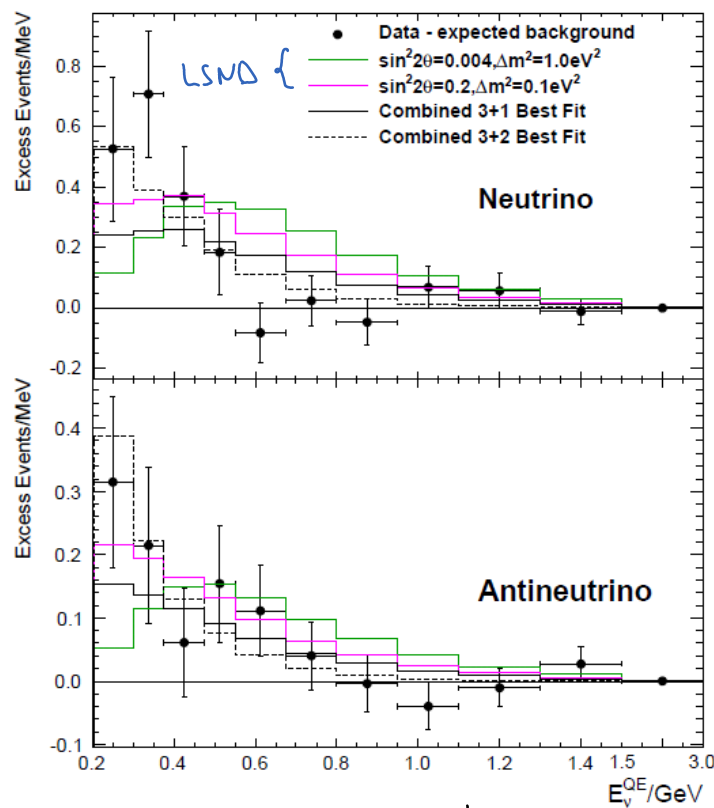
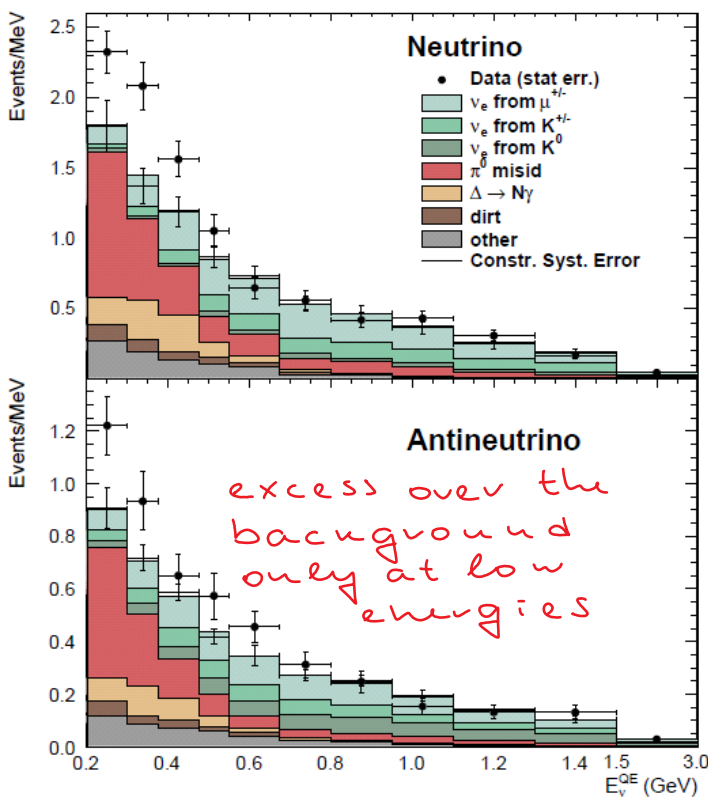
- the detector is located **541 m** away from the decay pipe.

- taking average neutrino energy  $E_\nu \sim 2 \text{ GeV}$  we find for  $E/L$ :

$$\frac{2 \times 10^3 \times 1 \text{ MeV}}{0.5 \times 10^{-2} \times 100 \text{ km}} = 4 \times 10^5 \times 10^{-5} \text{ eV}^2 \sim \mathbf{4 \text{ eV}^2}$$

i.e. MiniBooNE is sensitive to the same range of  $\Delta m^2$  as LSND, but the different range of energies allows to obtain an independent cross-check of the LSND results.

• Results:



in both neutrino and antineutrino modes

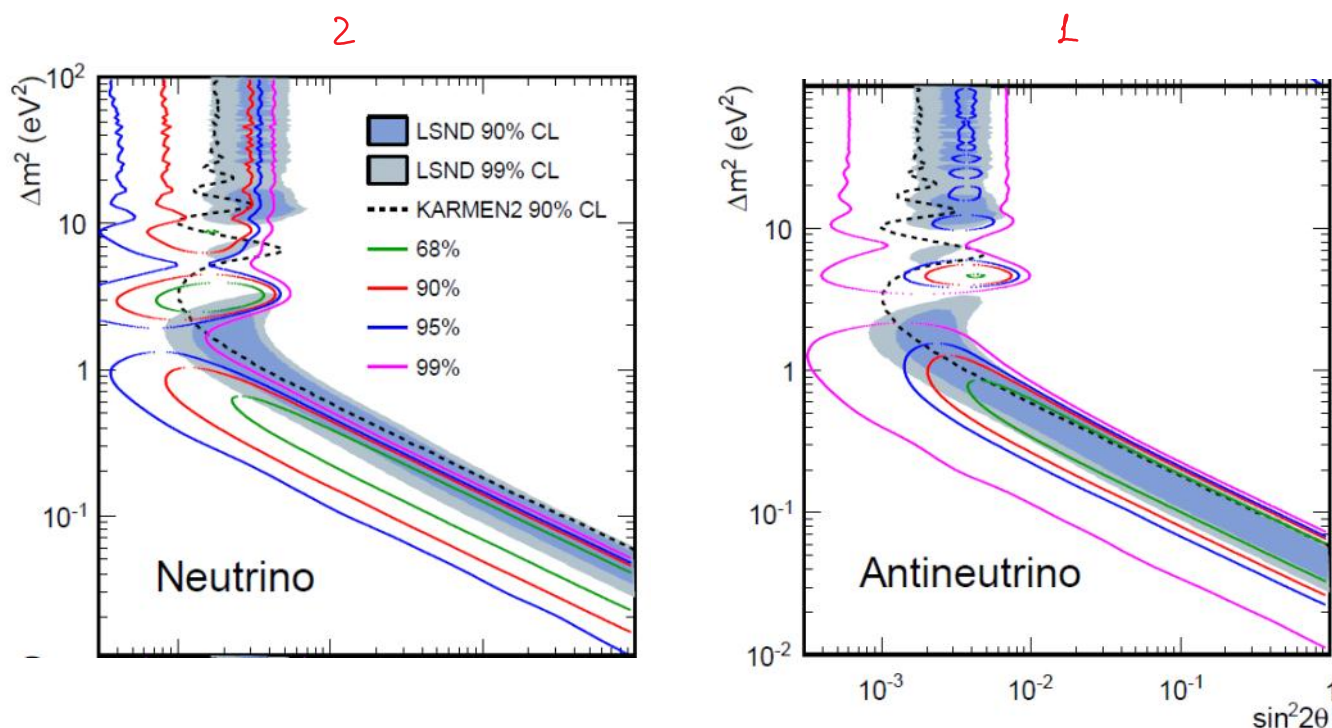
- At high energies no excess observed. At low

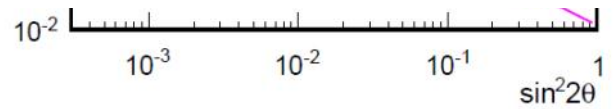
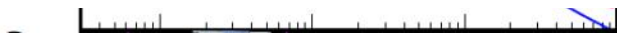
- At high energies no excess observed. At low energies there is an excess, but on the other hand, this is the region where the theoretical estimates are not particularly reliable.

- in the antineutrino mode the low-energy excess seems to be consistent with the LSND fit. But if we look more carefully, its shape is not consistent with simple 2-neutrino oscillation picture.

- in the neutrino mode the excess is larger and is worse compatible with the expectations based on the LSND results.

- interpretation of the neutrino and antineutrino results.





The antineutrino mode seems to be compatible with the results of LNSD, whereas in the neutrino mode the agreement is rather poor.

- This inconsistency can in principle be explained in 3+2 scheme, which introduces CP-violation and thus different properties of neutrinos and antineutrinos.

## Gallium anomaly

- Gallex and SAGE studied solar neutrinos using reaction  $\nu_e + {}^{71}\text{Ga} \rightarrow {}^{71}\text{Ge} + e^-$ .
- To calibrate the experiments two calibrations have been performed with  ${}^{51}\text{Cr}$  and  ${}^{37}\text{Ar}$  sources.
- They both have observed lack of events from the source, which cannot be explained by solar-type oscillations (orders of magnitude smaller distance).

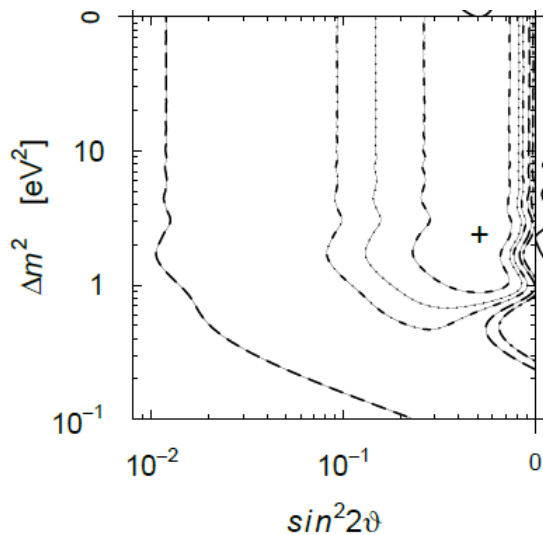
|                   | GALLEX                   |                          | SAGE                     |                               |
|-------------------|--------------------------|--------------------------|--------------------------|-------------------------------|
| k source          | G1<br>${}^{51}\text{Cr}$ | G2<br>${}^{51}\text{Cr}$ | S1<br>${}^{51}\text{Cr}$ | S2<br>${}^{37}\text{Ar}$      |
| $R_B^k$           | $0.953 \pm 0.11$         | $0.812^{+0.10}_{-0.11}$  | $0.95 \pm 0.12$          | $0.791 \pm^{+0.084}_{-0.078}$ |
| $R_H^k$           | $0.84^{+0.13}_{-0.12}$   | $0.71^{+0.12}_{-0.11}$   | $0.84^{+0.14}_{-0.13}$   | $0.70 \pm^{+0.10}_{-0.09}$    |
| radius [m]        |                          | 1.9                      |                          | 0.7                           |
| height [m]        |                          | 5.0                      |                          | 1.47                          |
| source height [m] | 2.7                      | 2.38                     |                          | 0.72                          |

- For Gallex the same detector, source and distance were used for the calibration and two different results were obtained.
- It is not quite clear, which numbers for the cross-sections are more reliable, but in any case the second measurement shows a deviation from the expected num.

in any case the second measurement shows a deviation from the expected number of events.

## - Interpretation.

Like for any other short-baseline experiment this result can be explained by the existence of an additional sterile neutrino. Implied parameters:

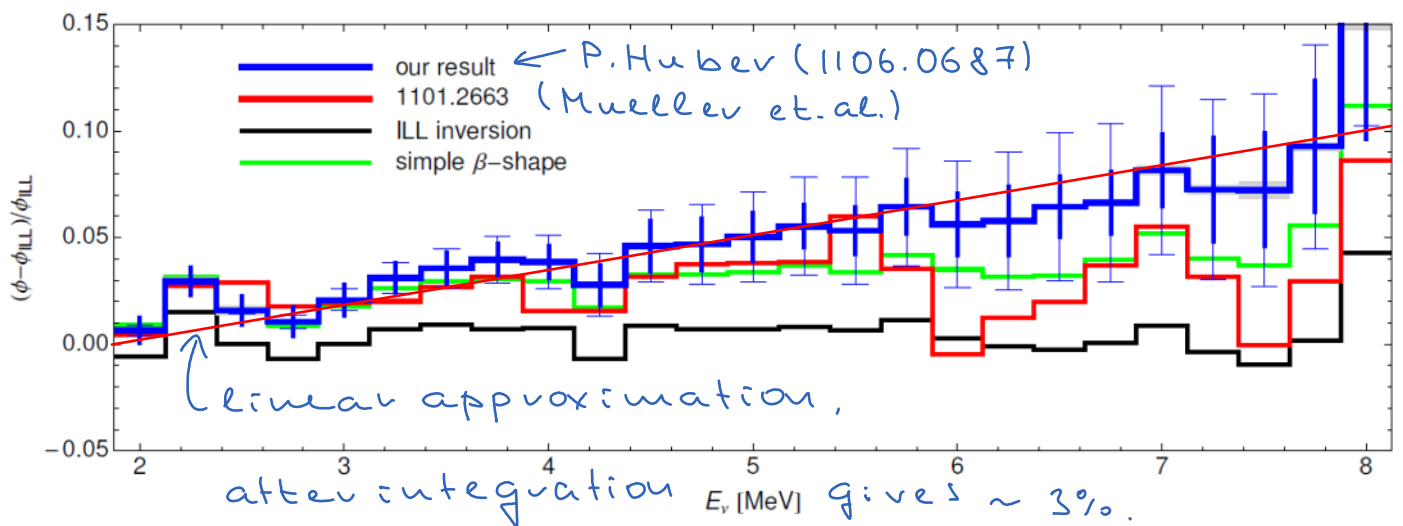


$$\sin^2 2\theta \gtrsim 0,07$$
$$\Delta m^2 \gtrsim 0,35 \text{ eV}^2$$

- Note that solar neutrino parameters have been measured by a number of experiments. Since the solar parameters obtained by Gallium and Sage are compatible with the others, we can consider the solar measurement as calibration and the source experiments as short-baseline sterile neutrino experiment.

## Reactor anomaly

- Reactor experiments belong to the ones most commonly used in the neutrino physics.
- Reactors produce a pure  $\bar{\nu}_e$  beam with typical energy  $E_\nu \sim$  a few MeV. This energy is too low to produce neutrinos or taus. Therefore, a reactor experiment can only be a disappearance experiment.
- There is a new estimate of the antineutrino flux that predicts  $\sim 3\%$  higher flux. The difference comes from higher energies.

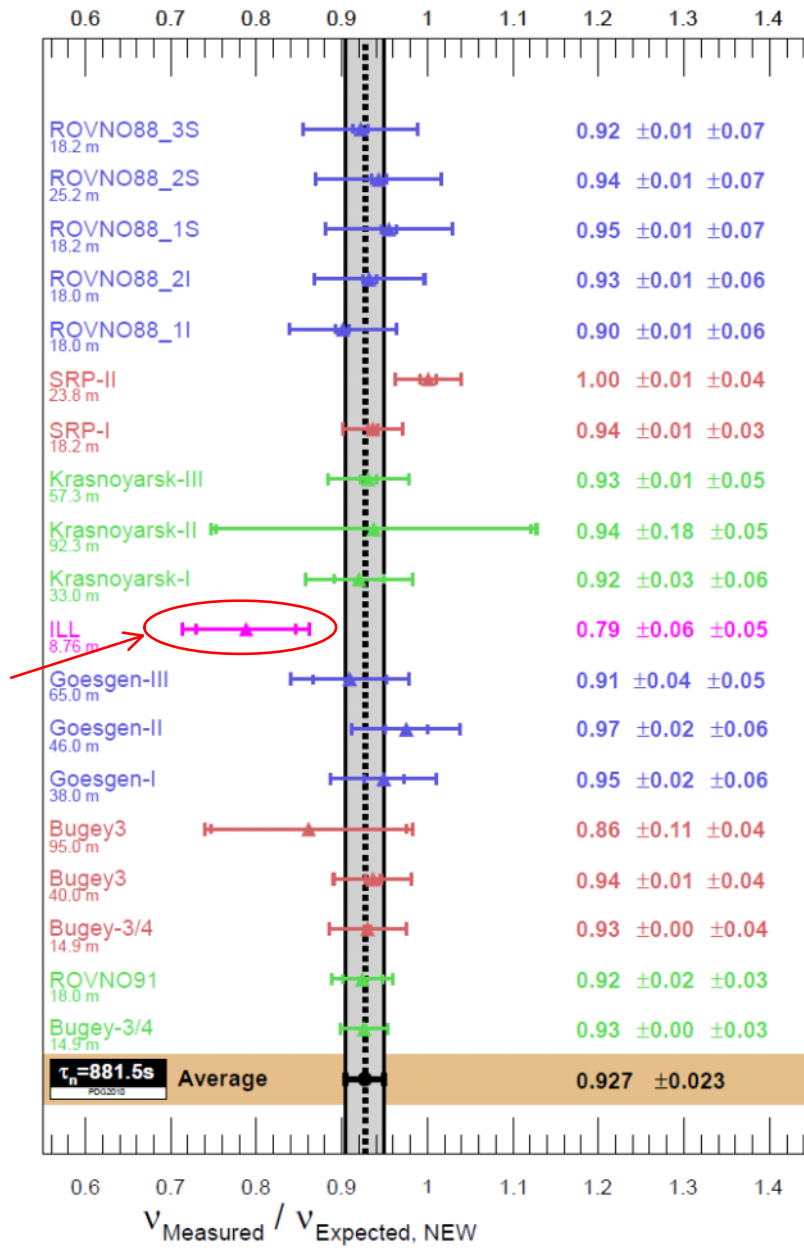


Note that a reliable estimate of the high-energy part of the spectrum is difficult. Therefore, even though there is no doubt that the new analysis improves the old results, it is unclear if the error-bars



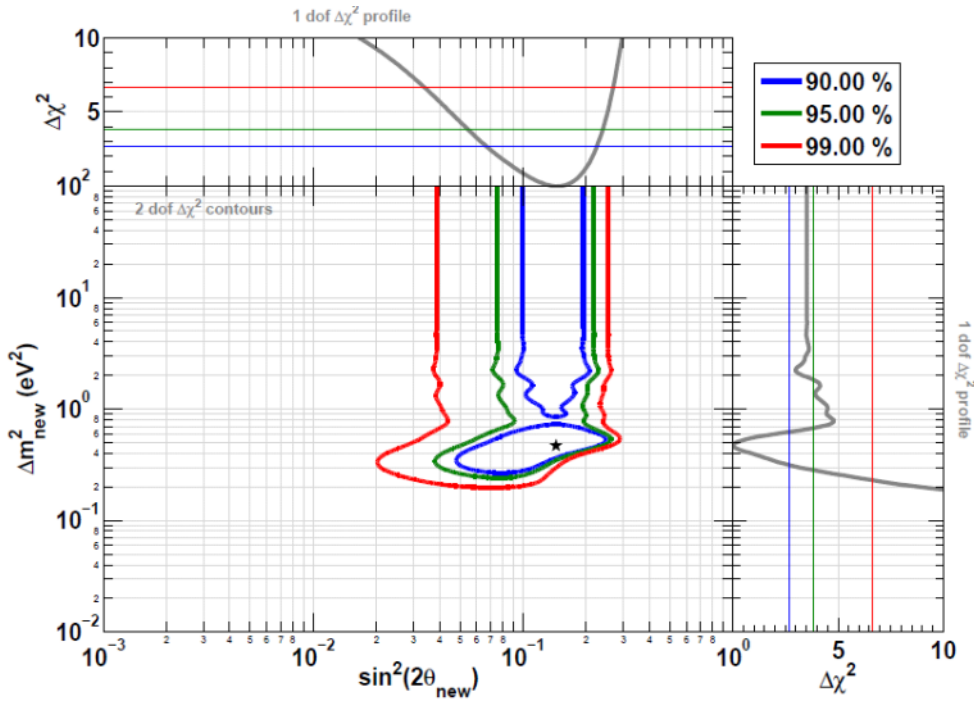
that the new analysis improves the old results, it is unclear if the error-bars are reliably estimated.

- Before the new estimate of the reactor spectrum, results of the old short-baseline oscillation experiments were compatible with the standard picture. But with the new theoretical fluxes they all have seen 3% too few events.
- Since typical distance from the reactor to the detector was below 100m, they were insensitive to  $\Delta m_{21}^2$  driven oscillations. Therefore, if we take the new prediction for the fluxes seriously, we can interpret the experimental results as a hint towards the existence of a new sterile neutrino.
- Reanalysis of the existing experimental data.



most of the experiments are consistent with each other.

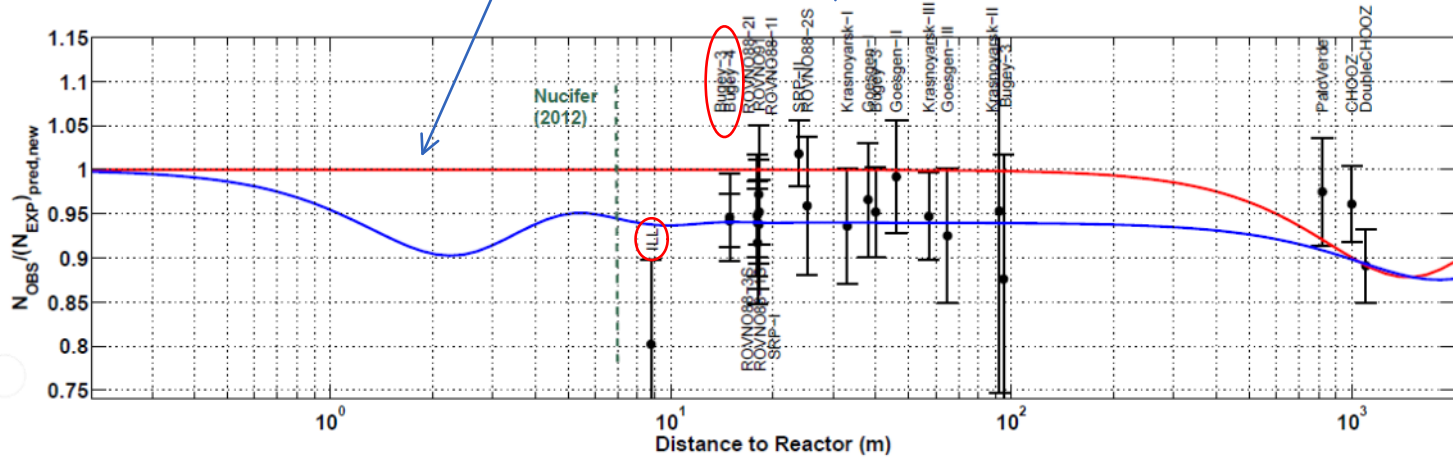
- Results of rate-only analysis:



vate-only:  
 $\Delta m^2 \approx 0.5 \text{ eV}^2$   
 $\sin^2(2\theta) = 0.15$

- The analysis based on the total flux alone strongly depends on the uncertainties in the flux calculation. It would be more interesting to observe L-dependence. This will probably be done by Nucifer experiment.

- Fit to data in 3 and 3+1 schemes. lies above most of the experiments



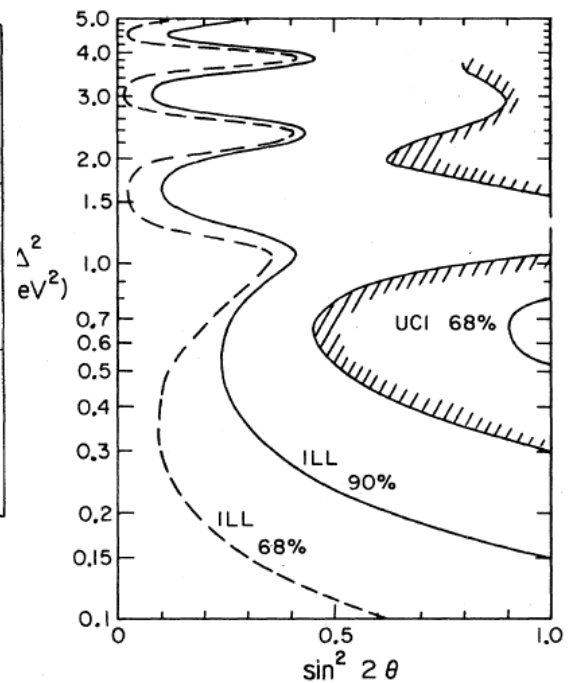
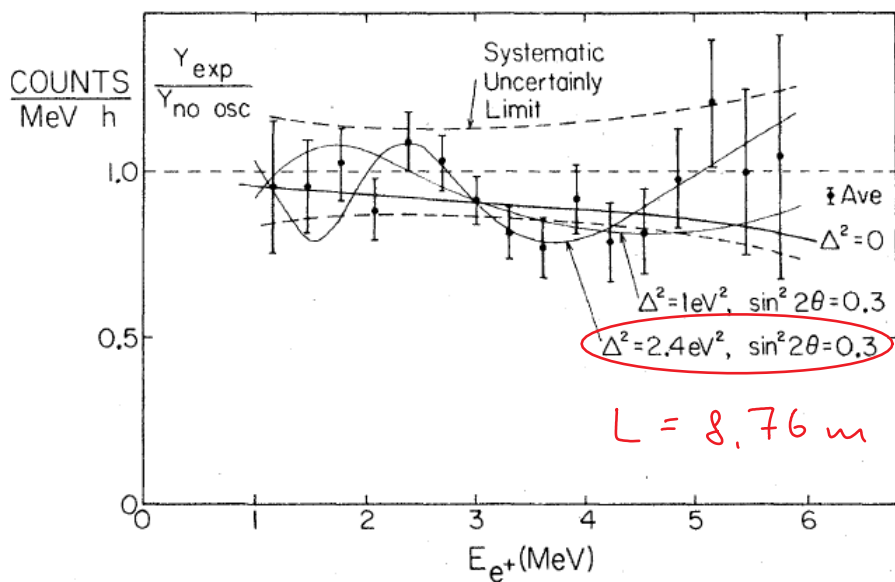
- red line assumes 3 neutrinos

- blue line assumes  $\sin^2(2\theta_{13}) \approx 0.085$  and a new neutrino with  $\Delta m^2 \gg 2 \text{ eV}^2$  and  $\sin^2(2\theta_{\text{new}}) = 0.12$ .

• Out of the past experiments only ILL and Bugey were close enough to the core to see the L-dependence.

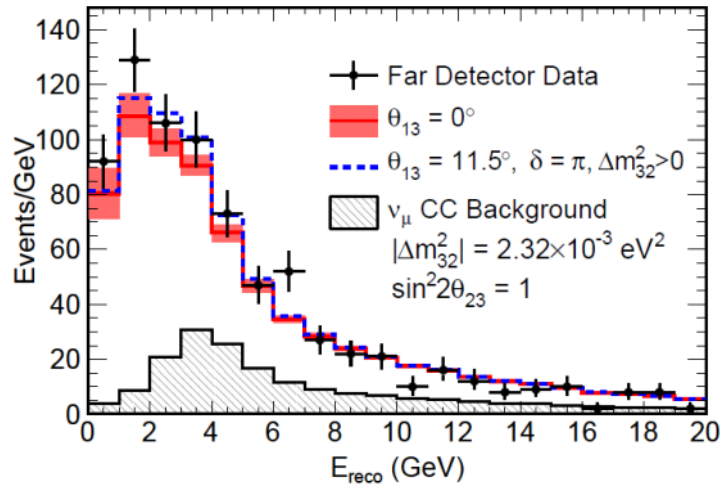
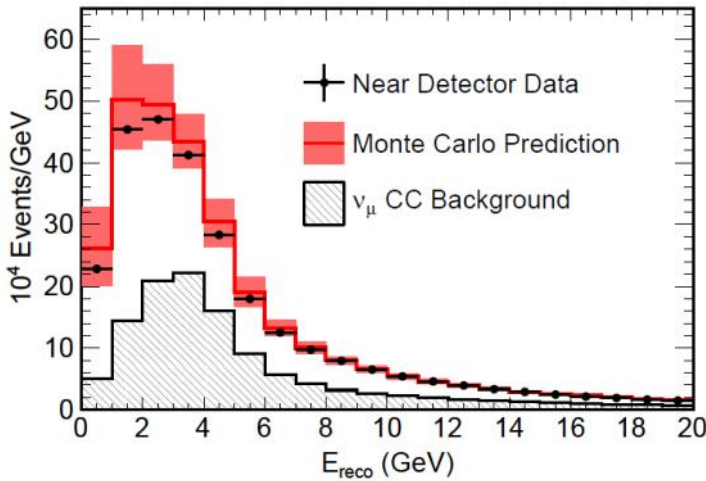
- Bugey didn't observe any distortion of the spectrum and therefore the new neutrino should be heavier than approximately  $1 \text{ eV}^2$ .

- On the other hand, there is an indication of the spectrum distortion in the ILL data. PRD 24, 1097 (1981)



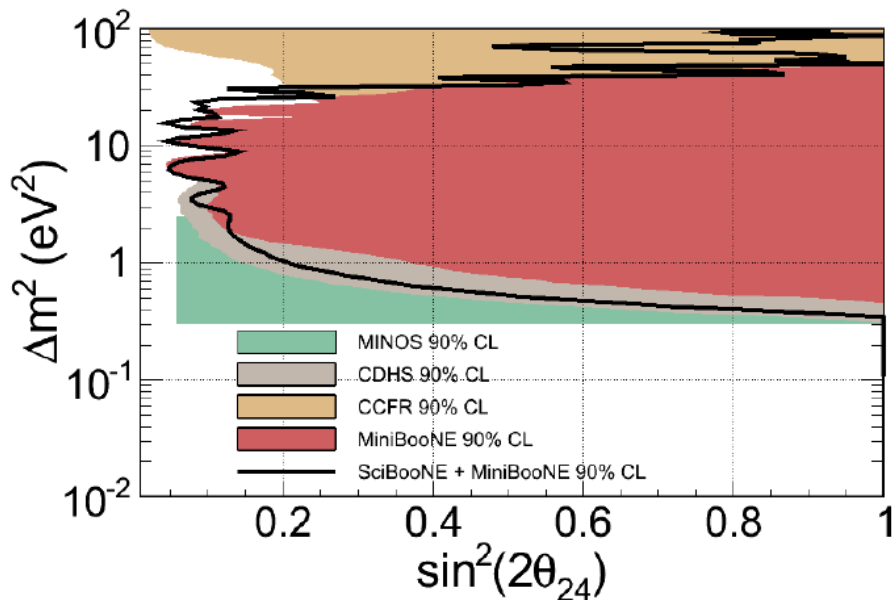


- Results: experimental results are compatible with 3V picture



- If we insist on the existence of the sterile neutrinos, then MINOS provides us with constraints on its parameters.

3+1  
scheme



- note that since the initial beam consists mainly of  $\nu_\mu$ , MINOS measures a different mixing angle as compared to the reactor experiments.

## Global fits

- To combine all the available data we need to specify a model. Data inconsistent within one model can be consistent within another.
- The anomalies are inconsistent with 3ν model. A possible minimal extension is a **3+1 model**.
- In this framework appearance and disappearance oscillation probabilities for **short baseline** experiments:

$$P_{\alpha \rightarrow \beta} = \sin^2(2\theta_{\alpha\beta}) \sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$$

$$P_{\alpha \rightarrow \alpha} = 1 - \sin^2(2\theta_{\alpha\alpha}) \sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$$

where

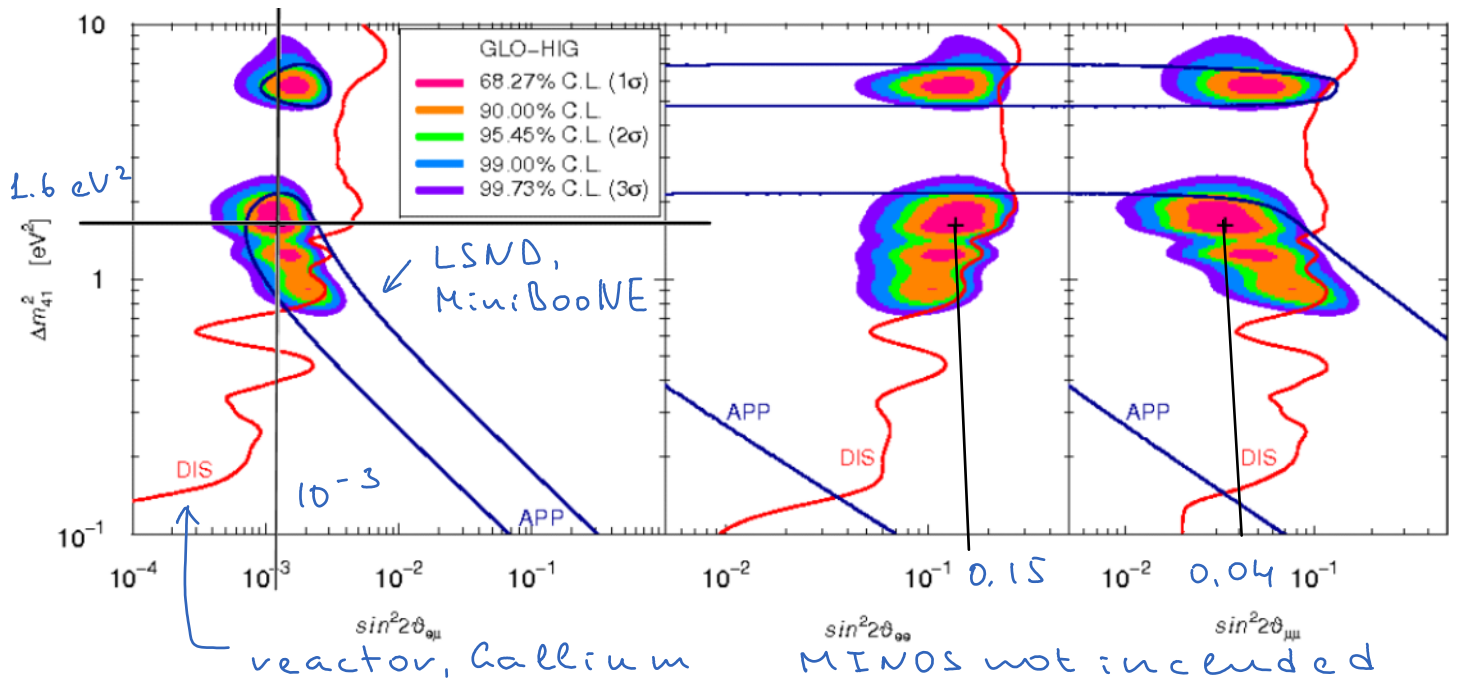
$$\sin^2(2\theta_{\alpha\beta}) = 4|U_{\alpha 4}|^2 |U_{\beta 4}|^2$$

$$\sin^2(2\theta_{\alpha\alpha}) = 4|U_{\alpha 4}|^2 (1 - |U_{\alpha 4}|^2)$$

- Theoretically, combining appearance and disappearance experiments we can measure the elements of the mixing matrix.

↙ Giunti et. al.





- since the low-energy data of MiniBooNE is somewhat controversial, I only show results of a global fit that does not include the low-energy data.

- There is a tension between the appearance and disappearance experiments. Especially for  $\theta_{13}$  mixing they are barely compatible.
- For comparison let me also show results of the global fit obtained by Thomas and Joachim:

SBL only: ← disappearance experiments similar to the fit of Giunti.

|     | $\Delta m_{41}^2$ [eV <sup>2</sup> ] | $ U_{e4} $ | $\Delta m_{51}^2$ [eV <sup>2</sup> ] | $ U_{e5} $ | $\chi^2/\text{dof}$ |
|-----|--------------------------------------|------------|--------------------------------------|------------|---------------------|
| 3+1 | 1.78                                 | 0.151      |                                      |            | 50.1/67             |
| 3+2 | 0.46                                 | 0.108      | 0.89                                 | 0.124      | 46.5/65             |

← appearance experiments  
 LSND and MiniBooNE data are in tension



~~up to date experiments~~  
 LSND and MiniBooNE data are in tension with the SBL data, as is evident from the next fit:

almost factor of four change.

|       | $ \Delta m_{41}^2 $ | $ U_{e4} $ | $ U_{\mu 4} $ | $ \Delta m_{51}^2 $ | $ U_{e5} $ | $ U_{\mu 5} $ | $\delta/\pi$ | $\chi^2/\text{dof}$ |
|-------|---------------------|------------|---------------|---------------------|------------|---------------|--------------|---------------------|
| (3+1) | (0.48)              | (0.14)     | (0.23)        |                     |            |               |              | (255.5/252)         |
| 3+2   | 1.10                | 0.14       | 0.11          | 0.82                | 0.13       | 0.12          | -0.31        | 245.2/247           |
| 1+3+1 | 0.48                | 0.13       | 0.12          | 0.90                | 0.15       | 0.15          | 0.62         | 241.6/247           |

## Summary

- To summarize: the experimental data are very inconclusive.
  - On one hand, there is a number of experiments which have observed a deviation from the standard picture.
  - On the other hand even experiments with very similar setups deliver different results. This makes one wonder if these results should be taken seriously.
  - Finally, results of the different experiments do not fit particularly well within simple theoretical frameworks. But this, of course, may be the problem of the theoretical frameworks themselves.